

# Graphene saturable absorber power scaling laser

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**Abstract:** A solution-processed graphene-film coated on a fiber-based connector is used for stable, mode-locked femtosecond-duration pulses with 16 mW average output power.

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## 1. Introduction

Graphene and carbon nanotubes (CNTs) are promising saturable absorbers (SAs) for mode-locking of fiber lasers [1–3]. Broadband operation is achieved with CNTs by combining tubes of different diameters [4], while it is an intrinsic property of graphene, due to the linear dispersion of Dirac electrons [5]. This, along with the ultrafast recovery time [6], and low saturation fluence [1, 7], makes graphene an excellent broadband SA [1, 7–9]. A variety of approaches have been used to make graphene-based SA (GSA) mode-locked lasers [10], with output power ranging, e.g., from  $\sim 3$  W [11] for a graphene oxide SA to  $\sim 270$  mW for a chemical vapor deposited GSA [12], generated in solid-state lasers [11, 12], to few-mW directly generated from fiber oscillators with GSAs coupled between fiber connectors [10]. The most commonly used GSA for fiber lasers [10] are fabricated from liquid phase exfoliation (LPE) of graphite [13, 14] in water [15] or organic solvents [13, 16]. Graphene produced by LPE can be embedded into polymer composites [2], which can be integrated into various systems [1, 2, 5, 7, 8]. Here we use LPE graphene in a polymer-free film coated onto a fiber based connector. This reduces non-saturable losses, making it suitable for high average-power applications and device miniaturization [8]. Based on this, we demonstrate a mode-locked fiber laser, achieving stable sub-250 fs pulses with a repetition rate of 21 MHz. The output power ranges from  $\sim 0.8$  mW to more than 15 mW.

## 2. Results

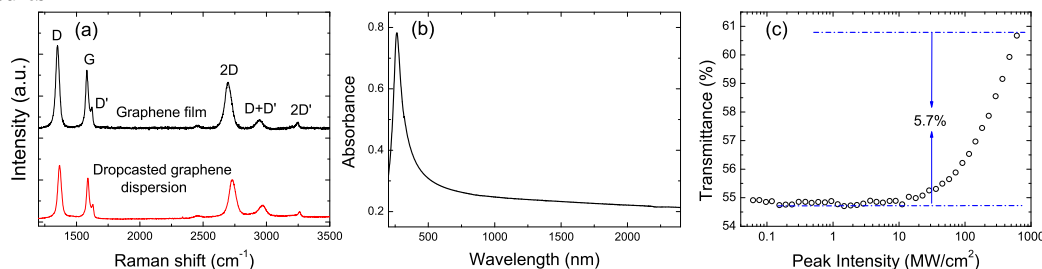


Fig. 1. (a) Raman spectrum of drop-cast-graphene and GSA. (b) Linear absorbance and (c) Nonlinear transmittance.

We fabricate a GSA by exfoliating graphite flakes via ultrasonic treatment in a solution of deionised water and sodium cholate (SC) (0.9wt%) [1, 15], followed by ultracentrifugation at 10000 rpm for 1 hour. The resulting top 70% dispersion is then filtered in vacuum through a nitrocellulose membrane (Millipore 100 nm pore-size filter). This blocks the flakes, while allowing water to pass through, resulting in a film on the top of the membrane. This is then transferred on the tip of a fiber connector and the membrane is dissolved in an acetone/water mixture. A control sample for optical characterization is also prepared by transferring the film on a quartz plate and annealing ( $\sim 90^\circ\text{C}$ , to improve adhesion) for 1 hour, followed by dissolution of the filter in acetone/water. To investigate the flakes' quality before and after film fabrication, and to monitor defects, we measure the Raman spectra at 457, 514.5, and 633 nm. Fig. 1(a) plots a typical Raman spectrum of a drop-cast graphene dispersion on a Si wafer and the resulting film. Besides the G and 2D peaks, significant D and D' bands as well as their combination mode D+D' at  $\sim 2950\text{ cm}^{-1}$  are also present [17]. We assign the D and D' peaks to the sub-micrometer edges of our flakes [18], rather than to disorder within the flakes. This is corroborated by a multi-wavelength Raman analysis: the resulting G peak dispersion is below  $0.05\text{ cm}^{-1}/\text{nm}$  [16]. There is no significant change in the spectrum of the film with respect to that of the drop-cast dispersion. Thus, the fabrication process does not affect the structure or quality of the flakes. The 2D peak is still single Lorentzian, thus,

even if the flakes are multi-layers, they are electronically decoupled and, to a first approximation, behave as a collection of single layers [17, 19]. The GSA has a featureless linear absorption from 500 to 2000 nm, Fig. 1(b), save the UV van Hove singularity peaks [20], and  $\sim 5.7\%$  non-linear transmittance change at  $1.5 \mu\text{m}$ , Fig. 1(c).

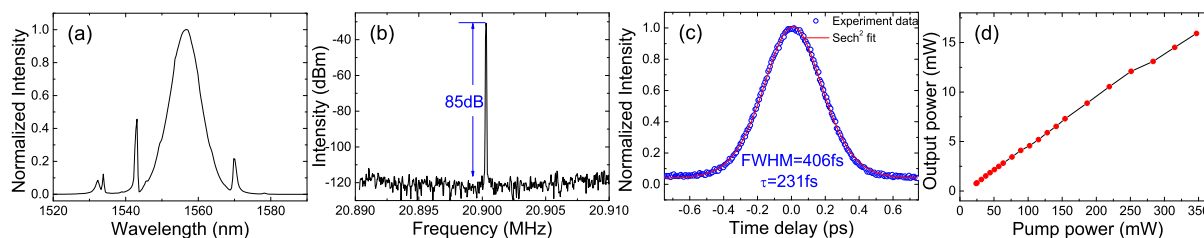


Fig. 2. (a) Optical spectrum; (b) RF spectrum (c); Autocorrelation trace and (d) Average output power vs. pump power.

We use a ring cavity with 3 m erbium doped fiber (EDF), pumped by a 980 nm diode through a wavelength division multiplexer. Unidirectional operation is enabled by an optical isolator. A polarization controller is used for mode-locking optimization. Output of the laser is provided by the 20% port of a coupler. The total cavity length is  $\sim 9$  m. Mode-locking starts at  $\sim 23$  mW pump power, with  $\sim 0.8$  mW output power at  $\sim 21$  MHz repetition rate. A typical spectrum is shown in Fig. 2(a), with the corresponding intensity autocorrelation trace in Fig. 2(c). The full width at half maximum (FWHM) is 406 fs. Assuming a  $\text{sech}^2$  fit, deconvolution gives 231 fs pulse duration. The radio frequency (RF) spectrum in Fig. 2(b) gives a signal-to-noise ratio  $> 80$  dB indicating low-amplitude fluctuations, thus stable mode-locking [21]. Mode-locked operation can be maintained up to a pump power  $\sim 350$  mW. The output power scales linearly with pump power as for Fig. 2(d), with a maximum  $\sim 16$  mW. The output power is limited by the maximum power of our pump laser, and higher average power would be possible using a higher power laser diode.

### 3. Conclusions

We demonstrated a mode-locked fiber laser using a graphene film saturable absorber. The easy integration of the GSA into the fiber connector shows potential in the development of next generation compact ultrafast fiber lasers.

### 4. Acknowledgments

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