

Raman spectroscopy of graphite in high magnetic fields: electron-phonon coupling and magnetophonon resonance

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Abstract. The magneto-Raman measurements of graphite were performed in a back-scattering Faraday geometry at temperature 10 K in magnetic fields up to 45 T. The experimental data reveal the rich structure of Raman-active excitations dominated by K-point massive electrons. At high magnetic fields the graphite E_{2g} Raman line shows complex multi- component behavior interpreted as magnetophonon resonance coupled electron-phonon modes at graphite's K-point. Also we found the clear signature of the fundamental, strongly dumped, $n=0$ magnetophonon resonance associated with H point massless holes.

Keywords: Raman scattering, magnetophonon resonance, graphite, high magnetic fields

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INTRODUCTION

The problem of electron-phonon coupling in graphene attracted recently significant attention. Electron-phonon coupling in graphene leads to the significant renormalization and broadening of long-wavelength E_{2g} optical phonons, what are predicted to occur through resonant interaction with Landau-quantized massless Dirac-like fermions in single layer graphene (SLG) [1,2] or massive fermions in bilayer graphene (BLG) [3]. In bulk graphite, the electron band structure combines the massless SLG-like holes and massive BLG-like electrons at the H and K points of the Brillouin zone, respectively. Manifestations of the magneto-phonon resonance (MPR) effect in SLG have been obtained in experiments of magneto-Raman scattering from the surface of bulk graphite and attributed to the presence of graphene flakes decoupled from the bulk graphite [4-8].

RESULTS AND DISCUSSION

In order to probe electron-phonon coupling and inter-Landau level electronic excitations in bulk graphite, we utilized the high magnetic field

Raman spectroscopy. The magneto-Raman measurements were performed in a back-scattering Faraday geometry at temperature 10 K in magnetic fields up to 45 T. [9] The experimental data reveal a rich structure of Raman-active excitations dominated by K-point massive electrons. At high magnetic fields the E_{2g} line shows complex multi- component behavior [9]. We interpret different observed modes as MPR coupled electron-phonon modes at graphite's K-point. Also we found a clear signature of the fundamental, strongly dumped, $n=0$ MPR associated with H point massless holes. In the highest field range ($B > 35$ T) the G line of bulk graphite narrows due to suppression of electron-phonon interaction.

Our results indicate on the necessity to use a full (Slonczewski –Weiss-McClure, SWM [10-12]) band structure model to describe high-field magneto-Raman results. It allowed us to extract accurately the values of SWM band parameters as well as electron-phonon coupling parameters at K- and H-points. In particular, the position and linewidth of unperturbed phonons can be derived from the high-field (> 35 T) spectra, i.e., the position of G-line, 1582.6 cm^{-1} , and the linewidth $\text{FWHM}_G = 4.4 \text{ cm}^{-1}$. We calculated the energies of asymmetric inter-Landau level (LL) transitions using SWM

model. In result, for coupling constants γ , g , and λ [13,14] we obtained $\gamma_r = 44 \pm 6 \text{ cm}^{-1}$, $g^{(K)} = 0.72 \pm 0.03 \text{ cm}^{-1}/\text{T}$, and extracted $\lambda_r^{(K)} \approx 3.3 \times 10^{-2}$, in agreement with the values of these parameters previously derived from density functional theory [14]. From analyses of the B - induced modulation of the central component of the G Raman peak near 32T [9], we concluded that this is consistent with MPR due to H -point inter-LL transitions, (1, 0) or (1, - 1), assuming that the LL widths are larger than the coupling strength. At the high-fields ($B > 30 \text{ T}$), the MPR-caused changes of G-line position fit well with the model when $\gamma_r = 100 \pm 10 \text{ cm}^{-1}$, $g_0^{(H)} \sim 3.2 \pm 0.2 \text{ cm}^{-1}/\text{T}^{1/2}$, and $\lambda_r^{(H)} \approx 1.6 \times 10^{-3}$. It is worth to point that $\lambda_r^{(H)}$ [9] is almost 20 times smaller than $\lambda_r^{(K)}$.

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